

# Light as a probe of fundamental physics in space

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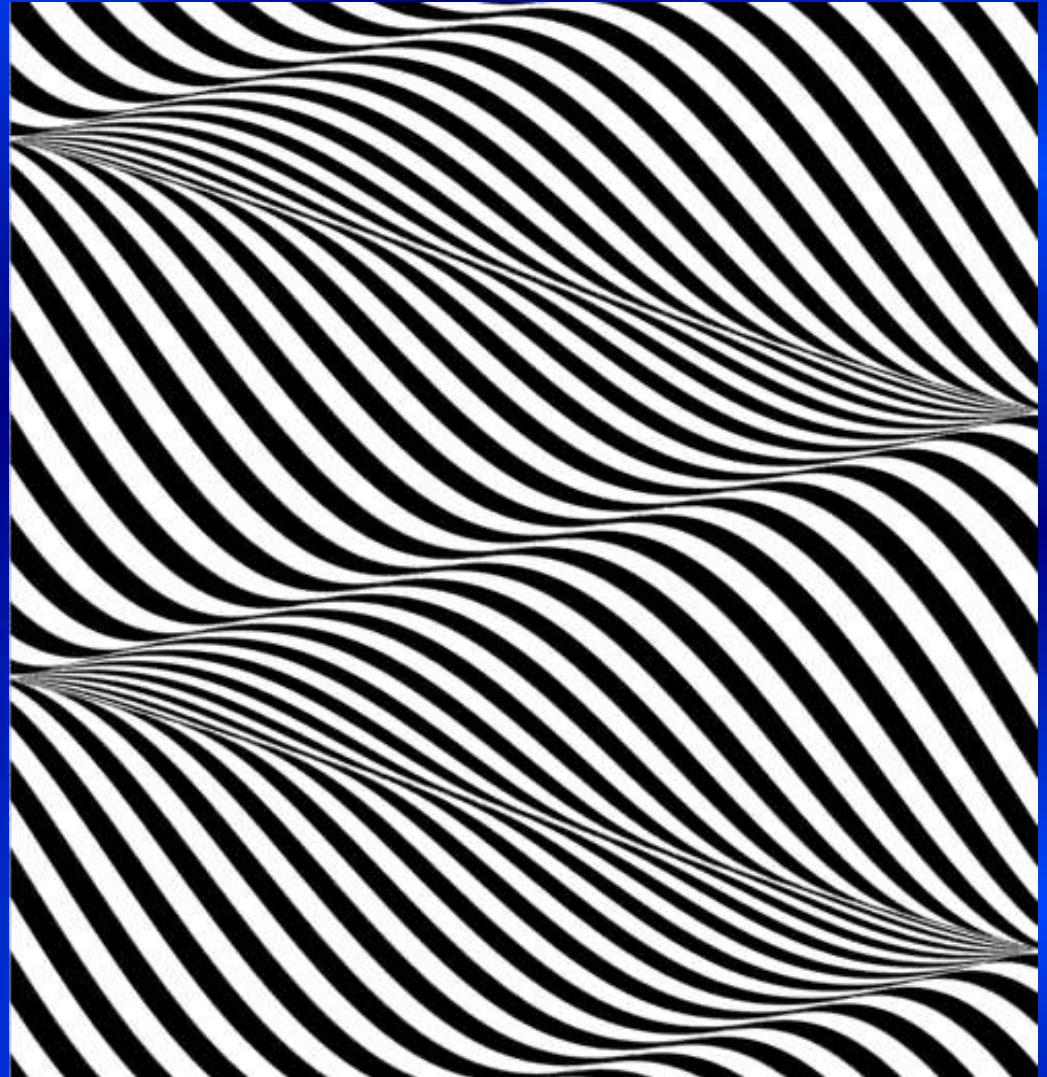
Politecnico di Torino and INFN

# Outline of the talk

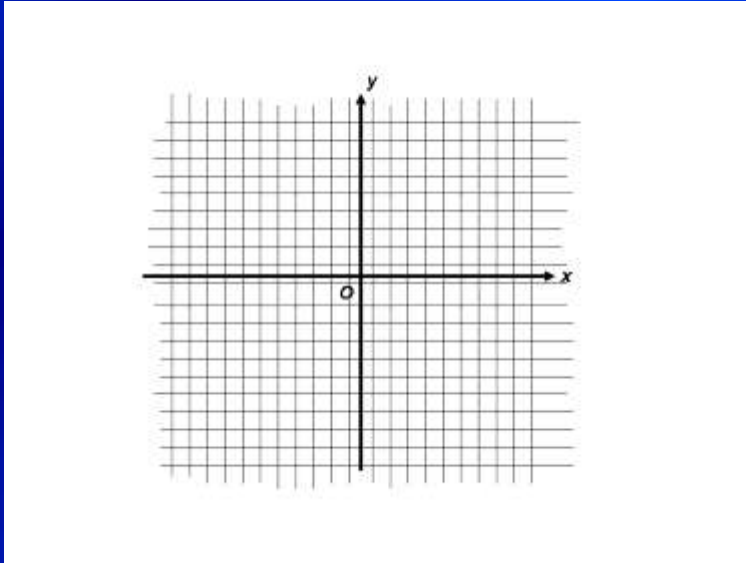
- Space-time and geodesics
- Null geodesics
- Space-time geodesy
- Light as a probe of the local structure of space-time
- Experiments in space

# Space-time and geodesics

Families of  
geodesics map  
any regular  
space- time  
patch

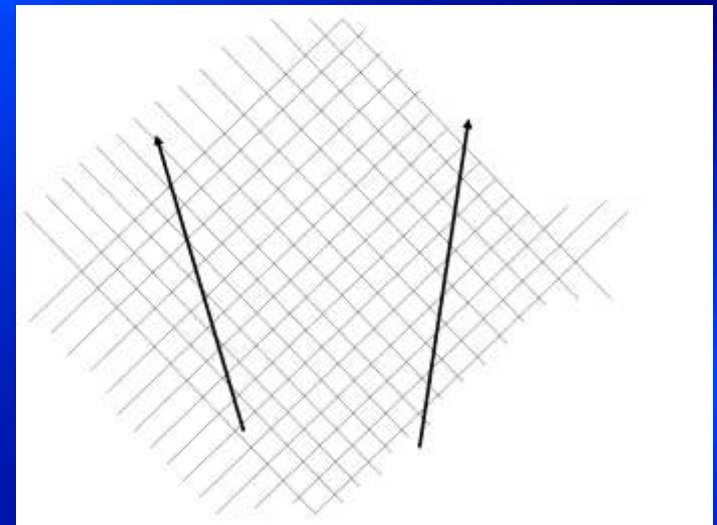


# Null or light coordinates



← Cartesian grid

Light rays grid →



# Wave vectors

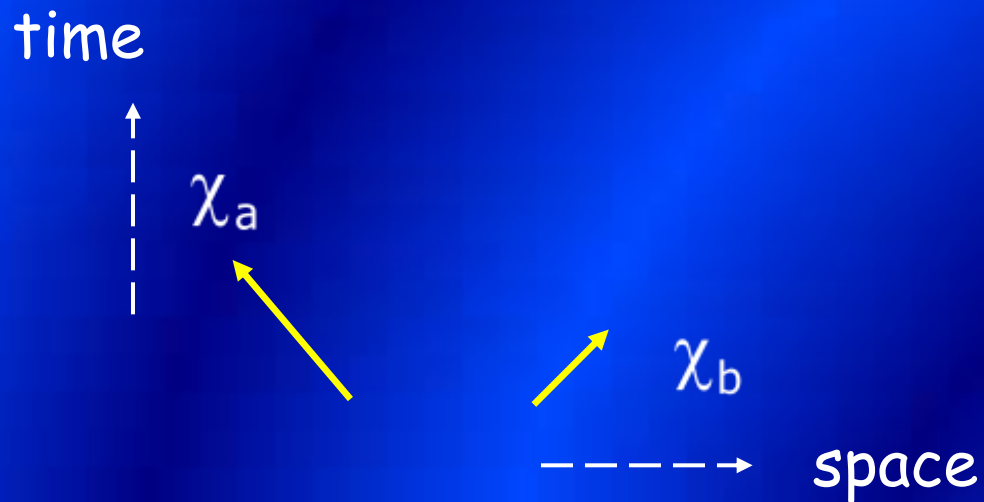
A null geodesic has a null tangent (wave) vector:

$$\chi = cT(1, \cos \alpha, \cos \beta, \cos \gamma) = cT(1, \hat{n})$$

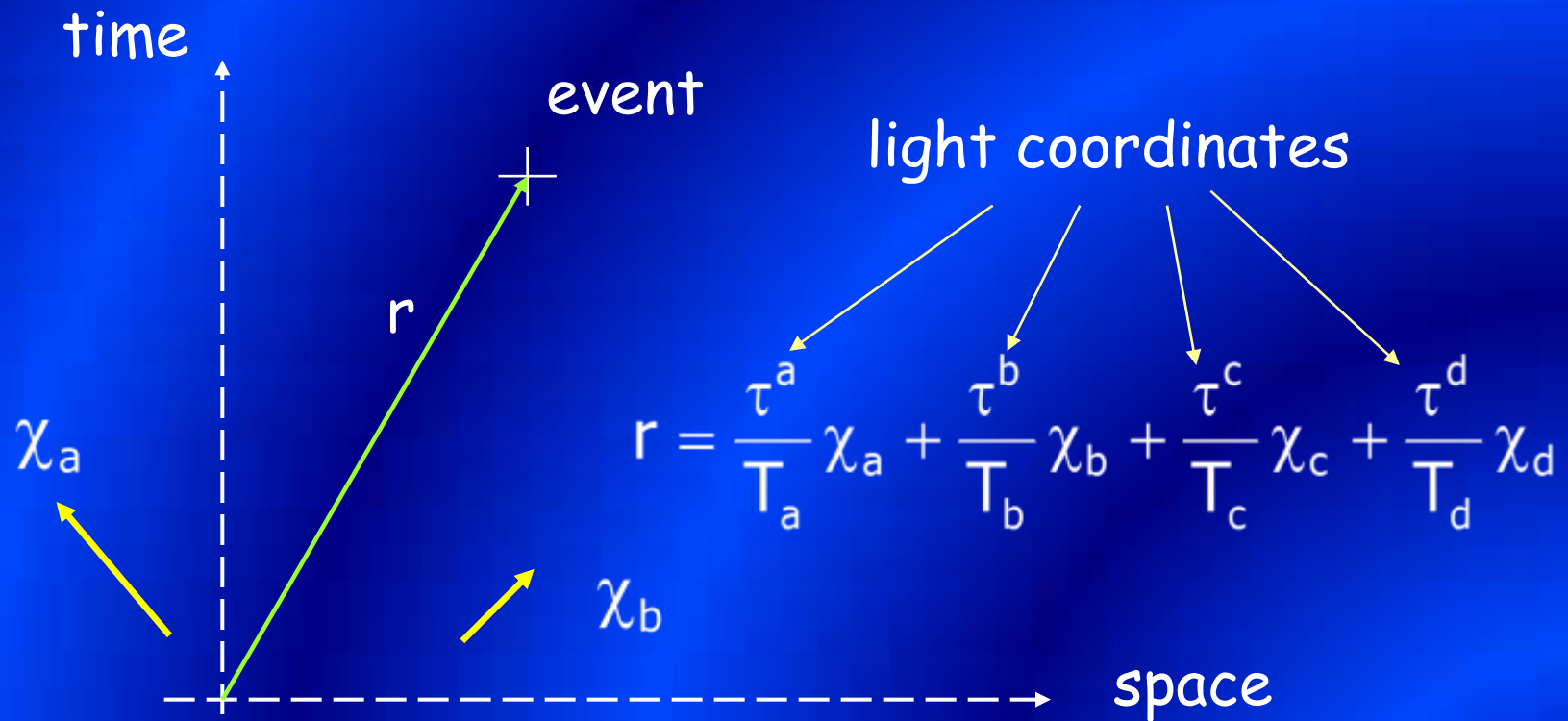
$$\chi^2 = 0$$

# Null bases

$\chi_a, \chi_b, \chi_c, \chi_d$



# Positioning in space-time



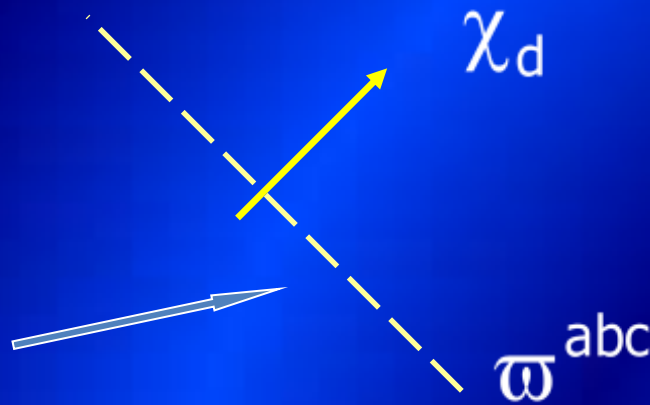
# Null wave fronts

Null  
hypersurface

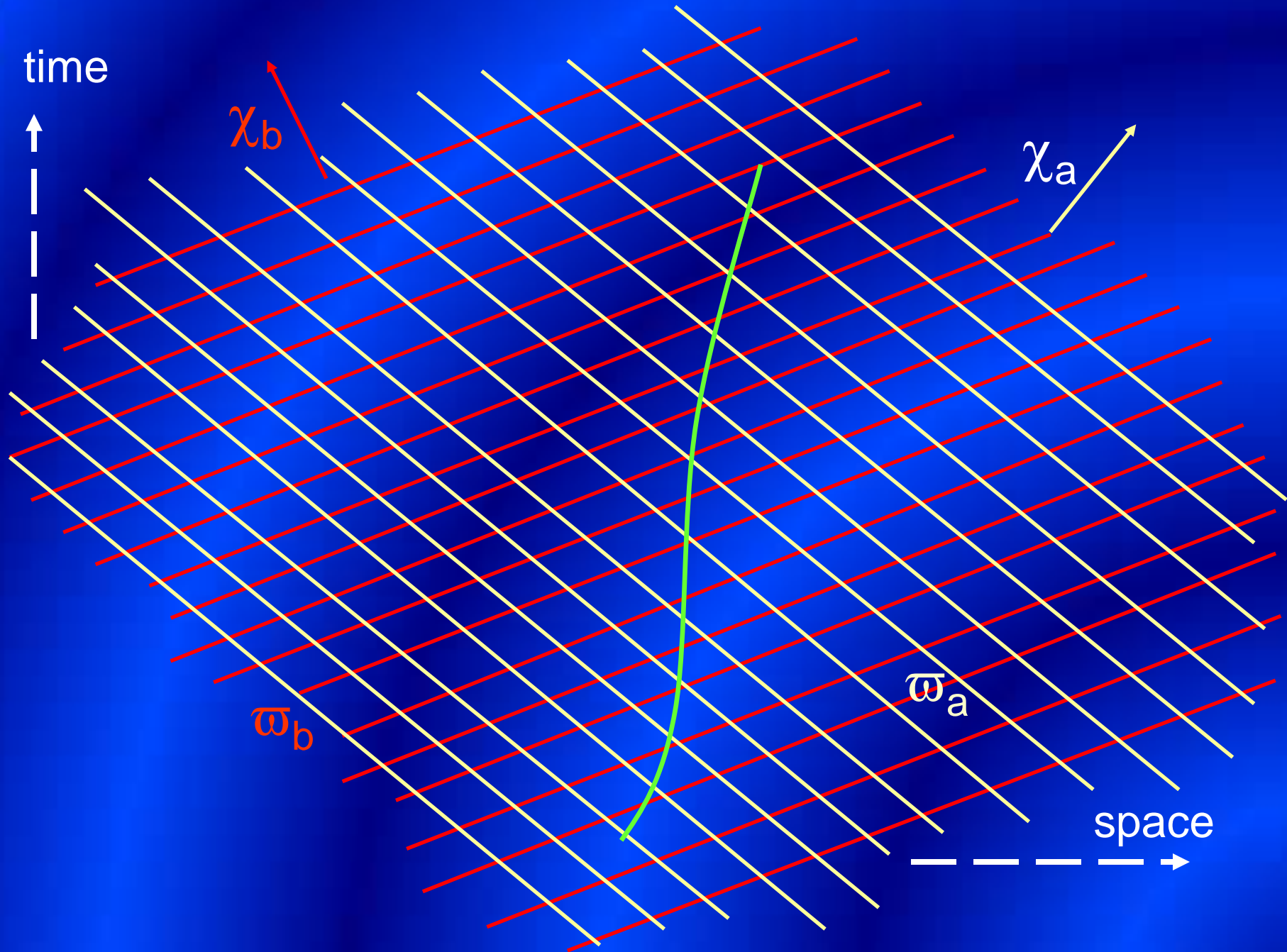


$$\varpi^{abc} = \varepsilon^{abcd} \chi_d$$

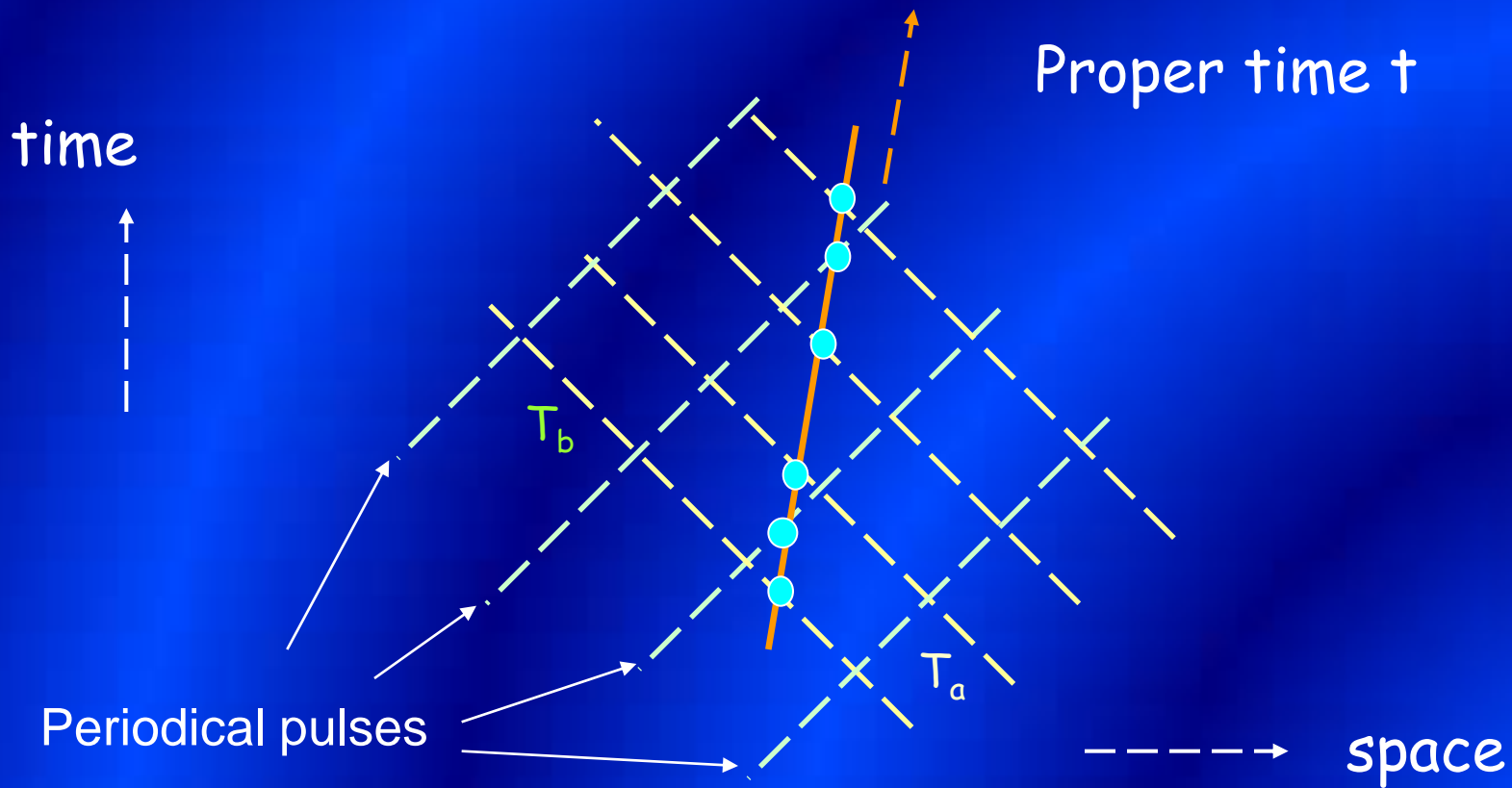
Locally a three-  
dimensional  
hyperplane







# Locally uniform motion



# Light coordinates of an event

$$\tau_{a,b,c,d} = [(n + x)\tau]_{a,b,c,d}$$

integer



From simple linear equations



# Uncertainty depends on clock

$$|\delta x| \leq 4 \left( \frac{1}{t_{i,i+4n}} + \frac{t_{i,i+1}}{t_{i,i+4n}^2} \right) \delta t$$

As big as allowed by the linearity of the worldline

# Accelerated motion

$$x^a = \frac{u^a}{T^a} t + \frac{1}{2} \frac{a^a}{T^a} t^2 + \dots$$

Four-velocity

Four-acceleration

Maximum integration time

$$t_{\max} = \sqrt{2 \left| \frac{u^a}{a^a} \right| \delta t}$$

# A gravitational field

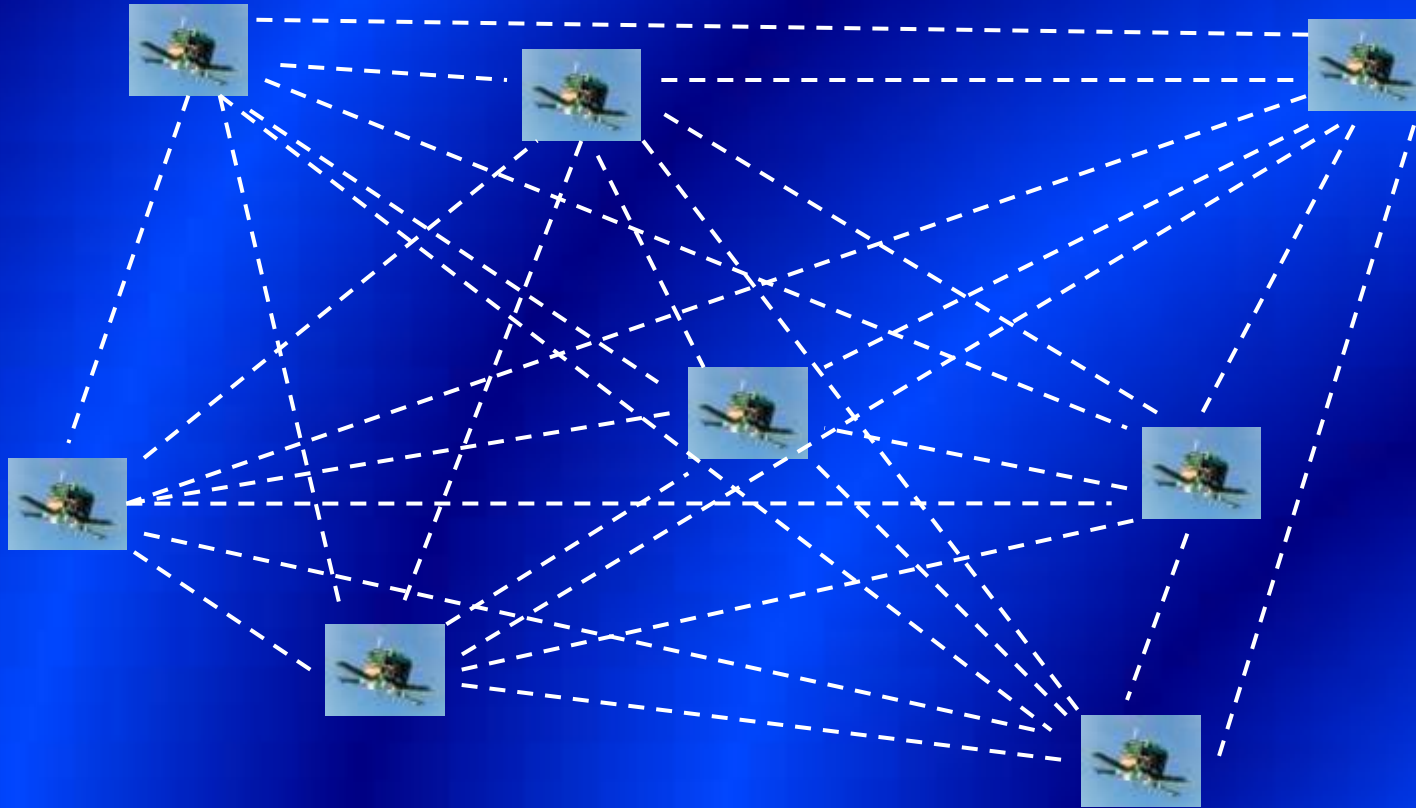
The gravitational field shows up when:

$$|\mathbf{u} \cdot \nabla \Phi| \geq 4 \frac{\delta t}{t^2}$$

Gravitational potential

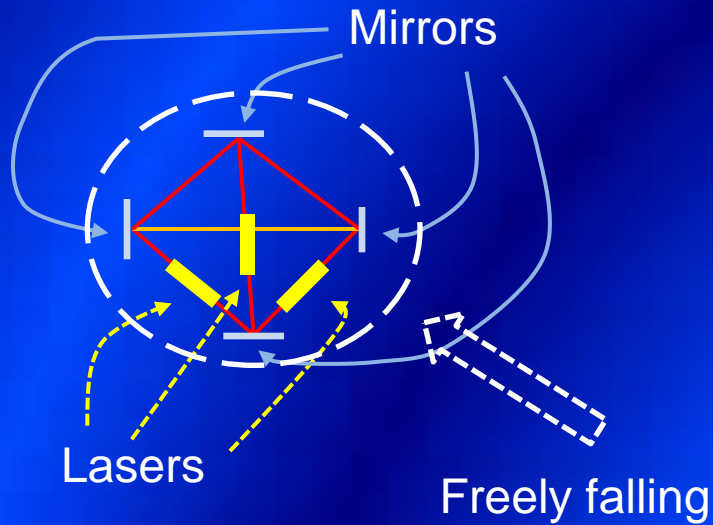


# Self positioning of a swarm of satellites (space-time geodesy)



Space-time geodesy

# Locally probing the curvature



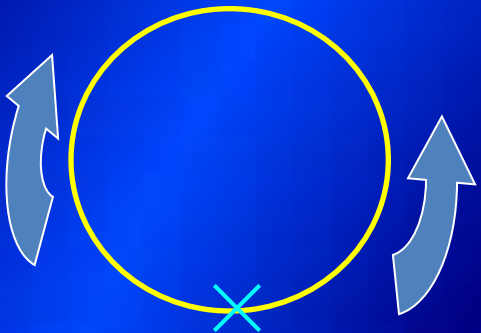
Tetrahedron with four  
triangular ringlasers





# The effect on frequency

Closed loop (shape irrelevant)



Counterrotating light beams

Time of flight difference

$$\delta T = T_+ - T_- = -2 \oint \frac{g_{0\phi}}{g_{00}} d\phi \neq 0$$

$$\delta \tau_0 = -2 \sqrt{g_{00}} \oint \frac{g_{0\phi}}{g_{00}} d\phi \neq 0$$

$$\Delta f = 4 \frac{S}{\lambda P} \left[ \frac{\omega}{c} \left( 1 + \frac{3 GM}{2 c^2 R} - \frac{3}{4} \sqrt{\frac{GM}{c^2 R}} + \frac{5 \omega^2 R^2}{2 c^2} \right) \hat{n}_a \cdot \hat{n}_S - \frac{1}{2} \frac{(GM)^{3/2}}{c^3 R^{5/2}} \hat{n}_\theta \cdot \hat{n}_S - \left( \frac{GJ}{c^3 R^3} \right) \hat{n}_g \cdot \hat{n}_S \right]$$

Beat frequency (equatorial orbit)

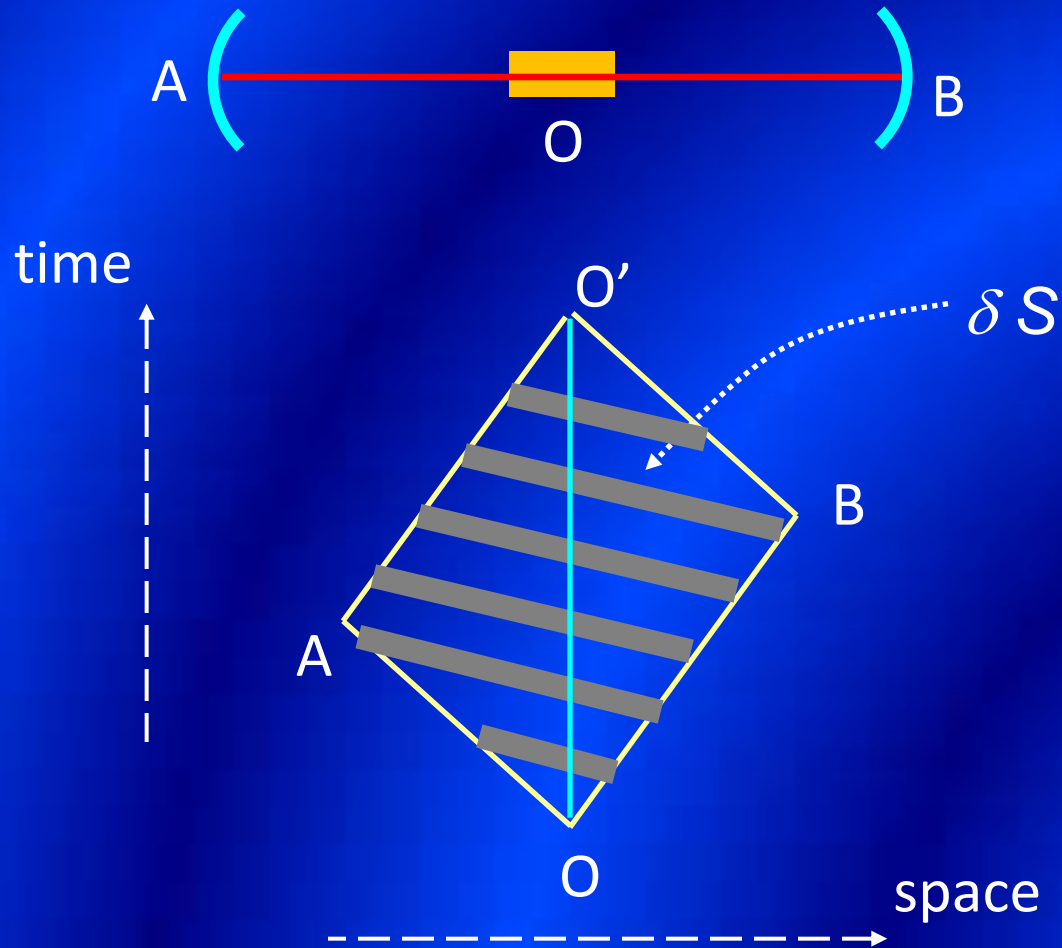
# Corotating device

$$\omega = \Omega = \sqrt{G \frac{M}{R^3}}$$

$$\Delta f = 4 \frac{S}{\lambda PR} \sqrt{\frac{GM}{R}} \left[ \left( 1 - \frac{3}{4} \sqrt{\frac{GM}{c^2 R}} + 4 \frac{GM}{c^2 R} \right) \hat{n}_a \cdot \hat{n}_S - \frac{1}{2} \frac{GM}{c^2 R} \hat{n}_\theta \cdot \hat{n}_S - \chi \left( \frac{\sqrt{GMR} \Omega_\oplus}{c^2} \right) \hat{n}_g \cdot \hat{n}_S \right]$$

$\sim 10^{-4}$        $\sim 10^{-9}$        $\sim 10^{-12}$

# A linear cavity



# The effect of curvature and anisotropy

$$\delta F^{\mu\nu} = \left( R^{\mu}_{\ \varepsilon 0 i} F^{\varepsilon\nu} + R^{\nu}_{\ \varepsilon 0 i} F^{\mu\varepsilon} \right) \delta S^{0i}$$

Riemann tensor

Electromagnetic tensor

Space-time area  
spanned by the cavity

Depends on the orientation of the cavity

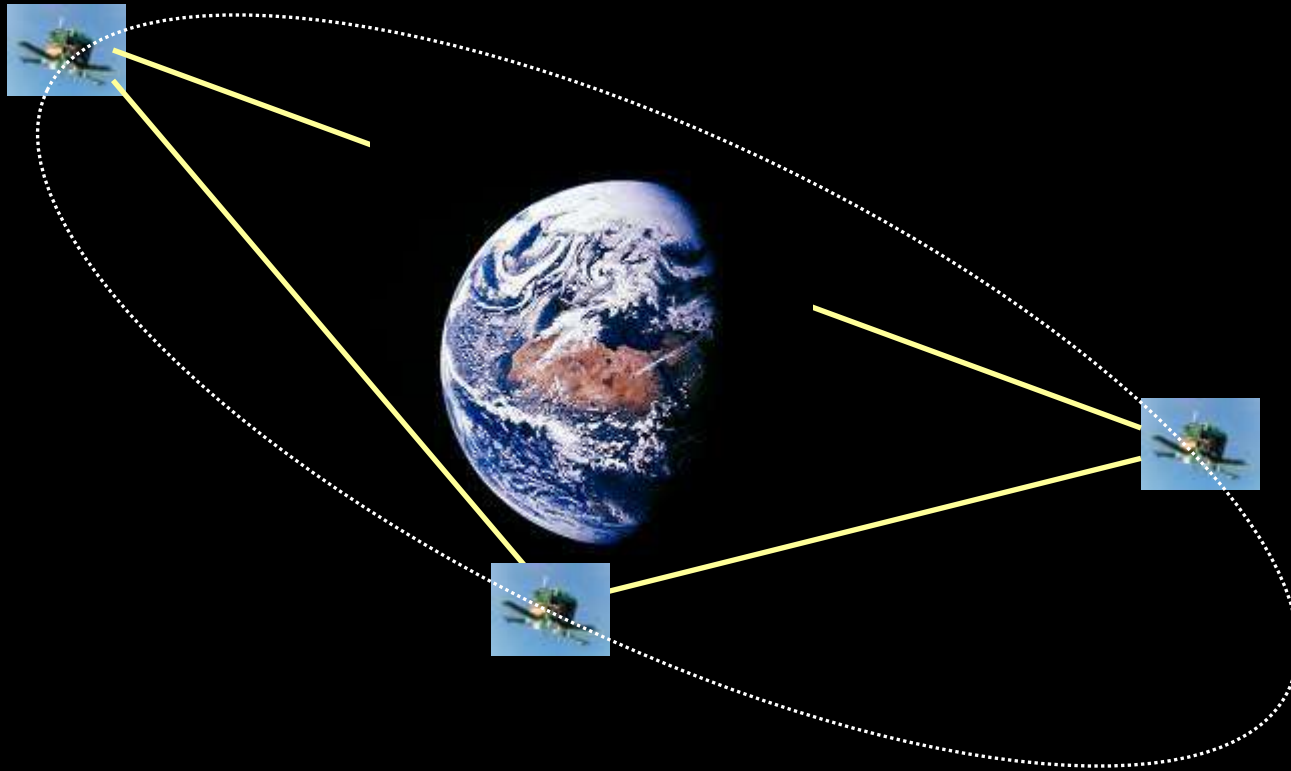
# Linearized Riemann tensor.

## An example:

$$R^{\phi}_{r0\theta} \approx \left[ \left( \frac{(GM)^{3/2}}{c^3 R^{7/2}} - 3 \frac{GJ}{c^3 R^4} \right) \frac{\cos \vartheta}{\sin^2 \vartheta} \right] \frac{l^2}{R} F^{\vartheta r}$$

$R^{\phi}_{r0\theta}$  → Length of the cavity  
 Change in the radial component of the magnetic field →  $\delta F^{\vartheta\phi}$   
 East-West component of the magnetic field →  $F^{\vartheta r}$

# Orbital gyro



# Arrival times difference

Equatorial orbit

$$\Delta t_0 \cong -6 \frac{R}{c} \left[ \left( 24 \sqrt{\frac{2}{\sqrt{3}}} + 3 \sqrt{\frac{\sqrt{3}}{2}} \right) \left( \frac{GM}{c^2 R} \right)^{\frac{3}{2}} - 2 \frac{\pi}{3} \frac{GM}{c^2 R} + 4 \frac{GJ}{c^3 R^2} + \sqrt{2\sqrt{3}} \sqrt{\left( \frac{GM}{c^2 R} \right)} \right]$$

Pulses

# Beat frequency

Equatorial orbit

$$\Delta f = \frac{c^2}{\lambda P} \Delta t_0 = \frac{c^2}{3\sqrt{3}\lambda R} \Delta t_0$$

Length of the loop





# PPN parameters

$$ds^2 = \left(1 - 2\frac{GM}{c^2 r}\right) d\tau^2 - \left(1 + 2\gamma\frac{GM}{c^2 r}\right) dr^2 - r^2 d\vartheta^2 \\ - r^2 \sin^2 \vartheta d\phi^2 + (4 + 4\gamma + \alpha_1) \frac{GJ}{2c^3 r} \sin^2 \vartheta d\tau d\phi$$

$$2M \rightarrow (1 + \gamma)M \qquad 2J \rightarrow \left(1 + \gamma + \frac{\alpha_1}{4}\right)J$$

# Conclusion

- Light is an intrinsically relativistic probe of space-time
- Closed paths of light in space (and space-time) evidence general relativistic effects either by
  - Interferometric techniques
  - or
  - Beat frequency measurements
- Electromagnetic pulses plus time of arrival times measurements outline a relativistic autonomous space navigation system.